

The matter's inertia and interaction in an isolated system

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Abstract: The intrinsic inertia and interaction of matters with unchanged density in an isolated system are studied. It is shown that the matters have a proved inertia of spin with the angular velocity $\vec{\omega}_n = \frac{d\theta}{dt} \hat{n} = \nabla \times \vec{u}$ of isolated system. The conclusive proof can reinforce Newton's the first law, taking into account the non-zero volume of point close to zero without limitation, and can explain the matter wave and seismic waves. A fundamental isolated system with two coupling matters has studied further. The revealed coupling characterizations have been used to explain the DNA structure, Time Cone and topological sphere of moving trace. The proven interaction within two matters is the coupling result of uniform rectilinear motion and spin of isolated system, which may be helpful to uniform the gravitation and electromagnetic force.

Introduction.—The conservation laws of matters can be classified as two distinct categories: those that are scalar including of mass, energy [1], electric charge; and those that are vector including of momentum, angular momentum, acceleration, angular acceleration, and spin of quantum particles as electron or atom [2,3]. Based on the viewpoint of the matters depending on ascertained space-time rather than mass point without any

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1 volume [4], the systems could be analyzed as followed.

2 *Model.*—A infinitesimal even isolated system is used as the physical
3 model, and its continuous equation of matters can be deduced using
4 Lagrangian method. According to the conservation law, the matters of
5 infinitesimal system keep constant as the product of density ρ and volume
6 V , $T=\rho V=C$, and its differential coefficient to the time parameter t is
7 listed as follows:

$$8 \quad \left. \frac{d\rho}{dt} \right|_V + \rho \left. \frac{dV}{Vdt} \right|_\rho = 0 \quad (1)$$

9 The centre of infinitesimal even isolated system is set up as the origin
10 of proprio-coordinate system to characterize its intrinsic properties,
11 similarly to the centre of mass. The velocity of any point in the system is
12 described by $\vec{u} = \frac{d\vec{l}}{dt}$, here $d\vec{l}$ for the displacement vector within time
13 quantum dt , and its divergence noted firstly in Chinese by Mohist Canon
14 is listed as follows,

$$\begin{aligned} 15 \quad \operatorname{div} \vec{u} &= \frac{1}{V} \lim_{V \rightarrow 0} \oiint_{\partial V} \vec{u} \cdot d\vec{S} = \frac{1}{V} \lim_{V \rightarrow 0} \oiint_{\partial V} \left(\frac{d\vec{l}}{dt} \right) \cdot d\vec{S} \\ 16 \quad &= \frac{1}{Vdt} \lim_{V \rightarrow 0} \oiint_{\partial V} d\vec{l} \cdot d\vec{S} \\ 17 \quad &= \frac{1}{Vdt} \lim_{V \rightarrow 0} \oiint_{\partial V} d(\vec{l} \cdot \vec{S}) = \frac{1}{Vdt} \lim_{V \rightarrow 0} V = \frac{dV}{Vdt} \\ 18 \quad \text{i. e.} \quad &\operatorname{div} \vec{u} = \frac{dV}{Vdt} \quad (2) \end{aligned}$$

19 Substitute equation (2) into equation (1), then

$$20 \quad \left. \frac{d\rho}{dt} \right|_V + \rho \operatorname{div} \vec{u}|_\rho = 0 \quad (3)$$

21 As the most important and fundamental theory during the transport

1 process in mechanics, the law of matter conservation with differential
 2 form is obtained, in which the density ρ or velocity u is variable for time t
 3 with constant volume V or constant density ρ , respectively.

4 The following analysis can be conducted from equations (1-3):

5 **Part (I):** *Kinematics of a single homogeneous matter within an*
 6 *isolated system.*— If there is a single homogeneous matter T_1 in the
 7 isolated system as shown in Dao De Jing, like deformable soft-body [5]
 8 water or rigid top, and its volume V_1 fills all space of the system [6].

$$\begin{aligned} 9 \quad T_1 &= \lim_{V_1 \rightarrow 0} \rho_1 V_1 = \rho_1 \lim_{V_1 \rightarrow 0} V_1 \\ 10 \quad \rho_1 &= \frac{T_1}{V_1} = \lim_{V_1 \rightarrow 0} \frac{T_1}{V_1} = \frac{dT_1}{dV_1} = C_{1(t)} \end{aligned}$$

11 From equation (3), then

$$12 \quad \left. \frac{d\rho_1}{dt} \right|_{V_1} = \left. \frac{d\left(\frac{dT_1}{dV_1}\right)}{dt} \right|_{V_1} = \left. \frac{d^2 T_1}{dV_1 dt} \right|_{V_1} = \left. \frac{dC_{1(t)}}{dt} \right|_{V_1} = -\rho_1 \operatorname{div} \vec{u}_1|_{\rho_1} \quad (4)$$

13 While $C_{1(t)}$ equals a non-zero and non-infinity constant to the physical
 14 meaning, $\rho_1 \operatorname{div} \vec{u}_1|_{\rho_1} = 0$ or $\operatorname{div} \vec{u}_1|_{\rho_1} = 0$ can be obtained from
 15 equation (4), and integrated further as

$$16 \quad \vec{u}_1 = \vec{C}' (C' \neq \pm\infty) \quad (5)$$

17 While the single homogeneous matter T of the system isn't fill space V
 18 fully, the rest space has the volume V' , density $\rho'=0$, matter $T'=0$ and
 19 velocity $\vec{u}' = 0$. So, $T=T_1+T'=T_1 \neq 0$, $V=V_1+V' \neq 0$ and from equation
 20 (3),

$$21 \quad \rho_1 = \frac{T_1}{V_1} = \lim_{V_1 \rightarrow 0} \frac{T_1}{V_1} = \frac{dT_1}{dV_1} = C_{1(t)}$$

$$\rho' = \frac{T'}{V'} = \lim_{V' \rightarrow 0} \frac{T'}{V'} = \frac{dT'}{dV'} = C'_{(t)} = 0$$

$$\bar{\rho} = \frac{T}{V} = \frac{T_1 + T'}{V} = \frac{\rho_1 V_1 + \rho' V'}{V} = \frac{C_{1(t)} V_1 + C'_{(t)} V'}{V} = \frac{V_1}{V} C_{1(t)} = \frac{V_1}{V} \rho_1$$

For the system,

$$\left. \frac{d\bar{\rho}}{dt} \right|_V + \bar{\rho} \operatorname{div} \bar{\vec{u}} \Big|_{\bar{\rho}} = 0$$

$$\left. \frac{dT}{dt} \right|_V + T \operatorname{div} \bar{\vec{u}} \Big|_{\bar{\rho}} = 0$$

$$\because \left. \frac{dT}{dt} \right|_V = 0 \quad \therefore \operatorname{div} \bar{\vec{u}} \Big|_{\bar{\rho}} = - \left. \frac{dT}{T dt} \right|_V = 0 \quad (4')$$

$$\text{Therefore} \quad \bar{\vec{u}} = \frac{\bar{u}_1 T_1 + \bar{u}' T'}{T_1 + T'} = \bar{\vec{C}}' = \bar{u}_1 (C' \neq \pm \infty) \quad (5)$$

Above equation revealed that the velocity of even matter keeps the same constant as the average velocity of the isolated system with single homogenous matter filled fully or not. If the matter is mass, the total momentum for the system keeps invariantly with the product of its mass and velocity.

Suppose that $\vec{l}_1 = l_1 \hat{\mathbf{r}}$ at the moment t , and in polar coordinate system with the same centre,

$$\begin{aligned} \bar{\vec{u}}_1 &= \frac{d\vec{l}_1}{dt} = \frac{dl_1}{dt} \hat{\mathbf{r}} + l_1 \frac{d\hat{\mathbf{r}}}{dt} = \frac{dl_1}{dt} \hat{\mathbf{r}} + l_1 \frac{d\theta}{dt} \hat{\mathbf{t}} = \frac{dl_1}{dt} \hat{\mathbf{r}} + \omega_n l_1 \hat{\mathbf{t}} \\ &= \bar{\vec{u}}_r + \bar{\vec{u}}_t + \bar{\vec{u}}_n = \bar{\vec{u}}_r + \bar{\vec{u}}_t \quad (\bar{\vec{u}}_n = 0 \hat{\mathbf{n}} = 0) \quad (6) \end{aligned}$$

Here the displacement vector $\hat{\mathbf{r}}$ is the unit vector for any point to the centre of system; vector $\hat{\mathbf{t}}$ is the unit vector perpendicular to the $\hat{\mathbf{r}}$ and lies in the plane consisted by the centre of system and unit vector $\hat{\mathbf{r}}$, $\hat{\mathbf{t}} \perp \hat{\mathbf{r}}$;

1 $\vec{\omega}_n = \frac{d\theta}{dt} \hat{n} = \nabla \times \vec{u}_1$ is the angular velocity of spin and its direction lines
 2 in unit vector \hat{n} which is perpendicular to unit vectors \hat{r} and \hat{t} ,
 3 simultaneously.

$$\begin{aligned}
 4 \quad |\vec{u}_1| &= \sqrt{\left(\frac{dl_1}{dt}\right)^2 + \left(l_1 \frac{d\theta}{dt}\right)^2} = |\vec{C}'| \\
 5 \quad &= \left| \sqrt{\left(\frac{dl_1}{dt}\right)^2 + \left(l_1 \frac{d\theta}{dt}\right)^2} \left[\frac{\frac{dl_1}{dt}}{\sqrt{\left(\frac{dl_1}{dt}\right)^2 + \left(l_1 \frac{d\theta}{dt}\right)^2}} \hat{r} \right. \right. \\
 6 \quad &\quad \left. \left. + \frac{l_1 \frac{d\theta}{dt}}{\sqrt{\left(\frac{dl_1}{dt}\right)^2 + \left(l_1 \frac{d\theta}{dt}\right)^2}} \hat{t} \right] \right|
 \end{aligned}$$

7 Given α as the angle between \vec{u}_1 and \hat{r} is equal to the product of ω_n
 8 and some parameter Δ , then

$$\begin{aligned}
 9 \quad \sin \alpha &= \frac{l_1 \frac{d\theta}{dt}}{\sqrt{\left(\frac{dl_1}{dt}\right)^2 + \left(l_1 \frac{d\theta}{dt}\right)^2}}, \quad \cos \alpha = \frac{\frac{dl_1}{dt}}{\sqrt{\left(\frac{dl_1}{dt}\right)^2 + \left(l_1 \frac{d\theta}{dt}\right)^2}} \\
 10 \quad \tan \alpha &= \frac{l_1 \frac{d\theta}{dt}}{\frac{dl_1}{dt}} = l_1 \frac{d\theta}{dl_1}, \quad \alpha = \tan^{-1} \left(l_1 \frac{d\theta}{dl_1} \right) = \omega_n \Delta
 \end{aligned}$$

$$\begin{aligned}
 11 \quad \text{Thus } \vec{u}_1 &= \sqrt{\left(\frac{dl_1}{dt}\right)^2 + \left(l_1 \frac{d\theta}{dt}\right)^2} [(\cos \alpha) \hat{r} + (\sin \alpha) \hat{t}] \\
 12 \quad &= A \omega_n [(\cos \omega_n \Delta) \hat{r} + (\sin \omega_n \Delta) \hat{t}] = \vec{C}' \\
 13 \quad (7)
 \end{aligned}$$

14 Here $A = \sqrt{\left(\frac{dl_1}{d\theta}\right)^2 + l_1^2}$, the velocity \vec{u}_1 can be identified as the wave with
 15 spin of ω_n in the space field.

16 *Discussions:*

1 1)If l_1 keeps constant, then $\alpha = 90^\circ$, $\vec{u}_1 \perp \hat{r}$, $\vec{u}_1 // \hat{t}$, the circling motion
 2 of matter can be observed on an axis lining with unit vector \hat{n} and
 3 crossing the centre of system. Considering the space volume of system
 4 is close to zero without limitation, $V \rightarrow 0$ and $V \neq 0$, the conclusion of
 5 the spin of matter can drawn with uniform angular velocity.

$$6 \quad \vec{u}_1 = \vec{u}_\tau = l_1 \frac{d\theta}{dt} \hat{t} = \omega_n l_1 \hat{t} = \vec{C}' \quad (8)$$

7 2) If θ keeps constant, then $\alpha = 0^\circ$, $\vec{u}_1 \perp \hat{t}$, $\vec{u}_1 // \hat{r}$, the uniform
 8 rectilinear motion of matter can be observed lining with unit vector \hat{r}
 9 and crossing the centre of system.

$$10 \quad \vec{u}_1 = \vec{u}_r = \frac{dl_1}{dt} \hat{r} = \vec{C}' \quad (9)$$

11 3) If l_1 and θ keep constant, then the matter of system keep static with
 12 $\vec{u}_1 = 0$. (10)

13 Supposing the matter is mass, the contents of equations (8), (9) and (10)
 14 are the proved and reinforced Newton's the first law [4], not for mass
 15 point with zero volume in space, but for the single homogeneous system
 16 with its non-zero volume close to zero with no limitation, and its density
 17 keeps a constant of neither zero nor infinity. The inertia of motive matter
 18 of an isolated homogeneous system including quantum particle, is to keep
 19 static state, uniform rectilinear motion or uniform angular spin [7].

20 4)If l_1 and θ are all changing, and α may be some ascertained value of
 21 0° to 90° , the matter performs the collective motive state of uniform
 22 rectilinear motion and uniform angular spin along the directions of \hat{r}

1 and $\vec{\tau}$, respectively, likening two types of seismic waves. As expressed in
 2 the equation of (7), the velocity \vec{u}_1 shows the motion of polarized
 3 transverse wave, where the properties of wave can be explained as the
 4 extension of uniform angular spin along the direction of uniform
 5 rectilinear motion, and its properties of particle can be explained as the
 6 uniform rectilinear motion, as the same as the matter wave in de
 7 Broglie hypothesis [8].

8 **Part (II): Kinematics and dynamics of two homogeneous matters**
 9 *within an isolated system.*—If two homogeneous matters T_1 and T_2 in an
 10 isolated system as noted in The Book of Changes, their volumes V_1 and
 11 V_2 are not filled all space of the system. At the initial time t_0 , an the origin
 12 O of proprio-coordinate system is signed by the centre of infinitesimal
 13 isolated system, and it will change to O_1 at time t_1 and further change to
 14 O_d at time t with a range of $dt=t-t_1$. As for the spins of matters T_1 and T_2
 15 can be ignored, considering that the supposed distances between matters
 16 and the centre of isolated system are very larger than their dimensions, all
 17 the movement about the integral system is discussed for the two-point
 18 model of matters T_1 and T_2 as shown in Dao De Jing. For the matter point
 19 of T_1 , its location can be described as a vector \vec{r}_{11} to the centre of O_1 and
 20 another vector \vec{r}_1 to the centre of O_d . The similar vectors \vec{r}_{21} and \vec{r}_2 can
 21 be listed for matter point of T_2 . Therefore, the points of centres O , O_1 and
 22 O_d can mark a plane of $\overline{OO_1O_d}$, $\vec{l} = \overline{OO_1}$ and $d\vec{l} = \overline{O_1O_d}$ are defined

1 as the displacement and unit one of the whole isolated system, as shown
2 in **Fig. 1**. The characters of matters T_1 , T_2 and rest space in the isolated
3 system are shown as $(T_1, V_1, \rho_1, \vec{r}_{11}, \vec{r}_1, \vec{u}_1; T_2, V_2, \rho_2, \vec{r}_{21}, \vec{r}_2, \vec{u}_2;$
4 $T_3=0, V_3, \rho_3=0, \vec{u}_3 = 0)$, and their relationships are
5 $T=T_1+T_2+T_3=T_1+T_2=C \neq 0$ and $V = V_1+V_2+V_3 \neq 0$.

6

7 FIG. 1. The geometric graph of moving matters T_1 and T_2 in an isolated system.

8 Therefore, from equation (3),

$$9 \quad \rho_1 = \frac{T_1}{V_1} = \lim_{V \rightarrow 0} \frac{T_1}{V_1} = \frac{dT_1}{dV_1} = C_{1(t)}$$

$$10 \quad \rho_2 = \frac{T_2}{V_2} = \lim_{V \rightarrow 0} \frac{T_2}{V_2} = \frac{dT_2}{dV_2} = C_{2(t)}$$

$$11 \quad \rho_3 = \frac{T_3}{V_3} = \lim_{V \rightarrow 0} \frac{T_3}{V_3} = \frac{dT_3}{dV_3} = C_{3(t)} = 0$$

$$12 \quad \bar{\rho} = \frac{T}{V} = \frac{T_1 + T_2 + T_3}{V} = \frac{\rho_1 V_1 + \rho_2 V_2 + \rho_3 V_3}{V} = \frac{C_{1(t)} V_1 + C_{2(t)} V_2}{V}$$

13 Put above parameters into the equation (3) for all isolated system

$$14 \quad \left. \frac{d\bar{\rho}}{dt} \right|_V + \bar{\rho} \text{div } \vec{\bar{u}} \Big|_{\rho} = 0$$

$$15 \quad \left. \frac{d(T_1 + T_2)}{dt} \right|_V + (T_1 + T_2) \text{div } \vec{\bar{u}} \Big|_{\rho} = 0$$

$$16 \quad (T_1 + T_2) \text{div } \vec{\bar{u}} \Big|_{\rho} = 0$$

$$17 \quad \text{That is} \quad \vec{\bar{u}} = \vec{C}' (C' \neq \pm \infty) \quad (5)$$

18 Above equation implied that the average velocity $\vec{\bar{u}}$ of two-matter
19 isolated system keeps constant vector \vec{C}' . If the matters are mass, the
20 total momentum \vec{P} of the system also keeps invariantly named the

1 conservation law of momentum as listed.

$$2 \quad \vec{P} = T_1 \vec{u}_1 + T_2 \vec{u}_2 = (T_1 + T_2) \vec{u} = (T_1 + T_2) \vec{C}' \quad (C' \neq \pm\infty) \quad (11)$$

3 And the average velocity of the whole isolated system is

$$4 \quad \vec{u} = \frac{T_1 \vec{u}_1 + T_2 \vec{u}_2}{T_1 + T_2} = \frac{T_1 \vec{u}_1}{T_1 + T_2} + \frac{T_2 \vec{u}_2}{T_1 + T_2} = \vec{C}'$$

$$5 \quad T_1 \vec{u}_1 + T_2 \vec{u}_2 = (T_1 + T_2) \vec{C}' \quad (C' \neq \pm\infty)$$

6 As the same as equation (6), here \vec{u} presents the collective motive state

7 of uniform rectilinear motion and angular spin of the whole system's

8 centre. Get derivative with respect to time, above equation comes into

$$9 \quad \frac{d(T_1 \vec{u}_1 + T_2 \vec{u}_2)}{dt} = T_1 \frac{d\vec{u}_1}{dt} + T_2 \frac{d\vec{u}_2}{dt} = \frac{d(T_1 + T_2) \vec{C}'}{dt} = (T_1 + T_2) \frac{d\vec{C}'}{dt}$$

10 Define $\vec{a}_1 = \frac{d\vec{u}_1}{dt}$ and $\vec{a}_2 = \frac{d\vec{u}_2}{dt}$, then

$$11 \quad T_1 \vec{a}_1 + T_2 \vec{a}_2 = 0$$

12 Further define $\vec{F}_1 = T_1 \vec{a}_1$ and $\vec{F}_2 = T_2 \vec{a}_2$, similarly as Newton's the

13 second Law [4], therefore $\vec{F}_1 + \vec{F}_2 = 0$ i.e. $\vec{F}_1 = -\vec{F}_2$ has obtained as

14 Newton's the third Law [4].

15 Based on the system centre O_1 , the respective proper orthogonal

16 decompositions of \vec{u}_1 , \vec{u}_2 , \vec{u} , and \vec{C}' vectors along \hat{r} , $\hat{\tau}$, and \hat{n} unit

17 vectors, can be substituted into Equation (11) as followed.

$$\begin{aligned} 18 \quad \vec{P} &= T_1 (\vec{u}_{1r} + \vec{u}_{1\tau} + \vec{u}_{1n}) + T_2 (\vec{u}_{2r} + \vec{u}_{2\tau} + \vec{u}_{2n}) \\ 19 \quad &= (T_1 + T_2) (\vec{u}_r + \vec{u}_\tau + \vec{u}_n) \\ 20 \quad &= (T_1 + T_2) (\vec{C}'_r + \vec{C}'_\tau + \vec{C}'_n) \quad (C' \neq \pm\infty) \quad (11') \end{aligned}$$

21 Therefore, the components of system along the \hat{r} , $\hat{\tau}$ and \hat{n} directions

1 are shown, following the same conservation law

$$2 \quad \vec{P}_r = T_1 \vec{u}_{1r} + T_2 \vec{u}_{2r} = (T_1 + T_2) \vec{u}_r = (T_1 + T_2) \vec{C}'_r \quad (11-1)$$

$$3 \quad \vec{P}_\tau = T_1 \vec{u}_{1\tau} + T_2 \vec{u}_{2\tau} = (T_1 + T_2) \vec{u}_\tau = (T_1 + T_2) \vec{C}'_\tau \quad (11-2)$$

$$4 \quad \vec{P}_n = T_1 \vec{u}_{1n} + T_2 \vec{u}_{2n} = (T_1 + T_2) \vec{u}_n = (T_1 + T_2) \vec{C}'_n \quad (11-3)$$

5 For \vec{u}_n is zero in equation (6), the equation (11-3) is transformed to

$$6 \quad \vec{P}_n = T_1 \vec{u}_{1n} + T_2 \vec{u}_{2n} = 0 \quad (11-3')$$

7 and further analyzed as, that the centre of two-matter isolated system has

8 no translation or rotation, but only harmonic vibration of two matters as

9 two spring oscillators around the centre and along direction \hat{n} for

10 astringency, their velocities follow the equation as

$$11 \quad \vec{u}_{1n} = -\frac{T_2}{T_1} \vec{u}_{2n} \quad (12-1)$$

12 and it performs the collective motive state of uniform rectilinear motion

13 and uniform angular spin in the plane of $\overline{OO_1O_d}$ consisted by \hat{r} and \hat{t}

14 unit vectors. Therefore, the translation, spin and vibration are orthogonal

15 along the \hat{r} , \hat{t} and \hat{n} directions, respectively.

16 And the similar equations to equation (12-1) can be deduced as

$$17 \quad \vec{u}_{1r} = -\frac{T_2}{T_1} \vec{u}_{2r} + \left(1 + \frac{T_2}{T_1} \vec{C}'_r\right) \quad (12-2)$$

$$18 \quad \vec{u}_{1\tau} = -\frac{T_2}{T_1} \vec{u}_{2\tau} + \left(1 + \frac{T_2}{T_1} \vec{C}'_\tau\right) \quad (12-3)$$

19 If the angles between directions of vector \hat{r} , \hat{t} or \hat{n} , and the line

20 crossing two matters and the centre of isolated system, are defined as α , β

21 or γ (α, β or $\gamma \in [0, \pi/2]$), the projection in direction \hat{r} of vector \vec{r}_1

22 is $\vec{r}_{1\perp} = \vec{r}_1 \cos \alpha$. Considering the same angular velocity $\vec{\omega}_n$ of spin and

1 the same line of \vec{r}_1 and \vec{r}_2 with the opposite direction for the centre of
2 two-mater isolated system, the directions of velocities $\vec{u}_{1\tau}$ and $\vec{u}_{2\tau}$ are
3 opposite, and their values fit the equation $\vec{u}_\tau = r_\perp \omega_n \hat{t}$ from $\vec{u}_\tau = \vec{\omega}_n \times$
4 \vec{r}_\perp , and the equation (11-2) is transformed to

$$\begin{aligned} 5 \quad \vec{P}_\tau &= T_1 r_{1\perp} \omega_n \hat{t} - T_2 r_{2\perp} \omega_n \hat{t} = (T_1 r_{1\perp} - T_2 r_{2\perp}) \omega_n \hat{t} \\ 6 \quad &= (T_1 r_1 \cos \alpha - T_2 r_2 \cos \alpha) \omega_n \hat{t} = (T_1 r_1 - T_2 r_2) \omega_n \cos \alpha \hat{t} \\ 7 \quad &= (T_1 + T_2) \vec{C}'_\tau = (T_1 + T_2) C'_\tau \hat{t} \end{aligned}$$

$$8 \quad \text{Therefore} \quad \omega_n = \frac{(T_1 + T_2) C'_\tau}{(T_1 r_1 - T_2 r_2) \cos \alpha} \quad (12-4)$$

9 *Discussions:*

10 1) When $\alpha = \pi/2$, $\vec{P}_\tau = 0$, $\vec{C}'_\tau = 0$, the equation (11-2) changes into the
11 same form as equation (11-3), and equation (12-4) is to be
12 meaningless. The collective motive states of polarizing motion with
13 contrary phases both in \hat{n} and \hat{t} directions, and of motionless or
14 uniform rectilinear motion in \hat{r} direction, should shape the moving
15 traces similar to the figures of Lissajous or twisted vines of plants [9].

16 2) When $\alpha \neq \pi/2$, $\omega_n = \frac{(\Delta T_1 + \Delta T_2) C'_\tau}{(\Delta T_1 r_1 - \Delta T_2 r_2) \cos \alpha}$, especially for $\alpha = 0$, then $\cos \alpha = 1$
17 and the minimum value $\omega_n = \frac{(\Delta T_1 + \Delta T_2) C'_\tau}{(\Delta T_1 r_1 - \Delta T_2 r_2)}$. The collective motive
18 states of polarizing motion with contrary phases in direction \hat{n} , of
19 spin with uniform the same angular velocity $\vec{\omega}_n$ in direction \hat{t} , and
20 of motionless or uniform rectilinear motion in direction \hat{r} , should
21 shape the moving traces similar to the familiar figures of double helix
22 like DNA structure which physical mechanism has not been revealed

before [10] or of funnel shape like Time Cone revealed by Einstein [11].

3) When $\vec{C}'_r = \vec{C}'_\tau = 0$, three equations of (11-1,2,3) are in the same formation returning to equation (11) after vector overlay operation. The collective motive states of polarizing motion with contrary phases in \hat{n} , $\hat{\tau}$ and \hat{r} directions should shape the moving trace of topological sphere, and the spins with contrary direction of two matters are both perpendicular to the surface of sphere [12].

4) Get derivative with respect to time, equations (12-1, 2, 3) become

$$\vec{a}_{1r} = -\frac{T_2}{T_1} \vec{a}_{2r} ; \vec{a}_{1\tau} = -\frac{T_2}{T_1} \vec{a}_{2\tau} ; \vec{a}_{1n} = -\frac{T_2}{T_1} \vec{a}_{2n} \quad (13)$$

The respective proper orthogonal decomposed equations (13) announce the relationship of $\vec{a}_1 = -\frac{T_2}{T_1} \vec{a}_2$ in the line of matters T_1 and T_2 , based on the centre O_1 of isolated system rather than that of matter T_1 or T_2 .

From equation (6), the velocities of matters T_1 and T_2 are listed as

$$\vec{u}_1 = \frac{d\vec{r}_1}{dt} = \frac{dr_1}{dt} \hat{r}_1 + r_1 \frac{d\hat{r}_1}{dt} = \frac{dr_1}{dt} \hat{r}_1 + r_1 \frac{d\theta}{dt} \hat{\tau}_1 = \frac{dr_1}{dt} \hat{r}_1 + \omega_n r_1 \hat{\tau}_1$$

$$= \vec{u}_{r1} + \vec{u}_{\tau1} + \vec{u}_{n1} = \vec{u}_{r1} + \vec{u}_{\tau1} \quad (\vec{u}_{n1} = 0 \hat{n}_1 = 0) \quad (6-1)$$

$$\vec{u}_2 = \frac{d\vec{r}_2}{dt} = \frac{dr_2}{dt} \hat{r}_2 + r_2 \frac{d\hat{r}_2}{dt} = \frac{dr_2}{dt} \hat{r}_2 + r_2 \frac{d\theta}{dt} \hat{\tau}_2 = \frac{dr_2}{dt} \hat{r}_2 + \omega_n r_2 \hat{\tau}_2$$

$$= \vec{u}_{r2} + \vec{u}_{\tau2} + \vec{u}_{n2} = \vec{u}_{r2} + \vec{u}_{\tau2} \quad (\vec{u}_{n2} = 0 \hat{n}_2 = 0) \quad (6-2)$$

Their accelerations are listed as

$$\vec{a}_1 = \frac{d\vec{u}_1}{dt} = \left[\frac{d^2 r_1}{dt^2} - r_1 \omega_n^2 \right] \hat{r}_1 + \left(2\omega_n \frac{dr_1}{dt} + r_1 \frac{d\omega_n}{dt} \right) \hat{\tau}_1 + 0 \hat{n}_1$$

$$\vec{a}_2 = \frac{d\vec{u}_2}{dt} = \left[\frac{d^2 r_2}{dt^2} - r_2 \omega_{n2}^2 \right] \hat{r}_2 + \left(2\omega_{n2} \frac{dr_2}{dt} + r_2 \frac{d\omega_{n2}}{dt} \right) \hat{t}_2 + 0\hat{n}_2$$

The relationship of unit vectors $(\hat{r}_1, \hat{t}_1, \hat{n}_1)$ and $(\hat{r}_2, \hat{t}_2, \hat{n}_2)$ has been listed as followed,

$$\begin{aligned}\hat{r}_2 &= -\hat{r}_1 \\ \hat{t}_2 &= -\hat{t}_1 \\ \hat{n}_2 &= \hat{n}_1\end{aligned}$$

Considering the isolated system has the same angular velocity with that of matters T_1 and T_2 , therefore, those angular velocities have the same value $\omega_{n1} = \omega_{n2} = \omega_n$ of binary system [13].

The forces \vec{F}_1 and \vec{F}_2 along and/or across the line between matters T_1 and T_2 are

$$\vec{F}_1 = T_1 \vec{a}_1 = T_1 \left[\frac{d^2 r_1}{dt^2} - r_1 \omega_{n1}^2 \right] \hat{r}_1 + T_1 \left(2\omega_{n1} \frac{dr_1}{dt} + r_1 \frac{d\omega_{n1}}{dt} \right) \hat{t}_1$$

$$\vec{F}_2 = T_2 \vec{a}_2 = T_2 \left[\frac{d^2 r_2}{dt^2} - r_2 \omega_{n2}^2 \right] \hat{r}_2 + T_2 \left(2\omega_{n2} \frac{dr_2}{dt} + r_2 \frac{d\omega_{n2}}{dt} \right) \hat{t}_2$$

With the formula $|\nabla \times \vec{u}|$ to instead of ω_n , the uniformity from the centre of isolated system is

$$\begin{aligned}\vec{F} &= T \vec{a} = T \left[\frac{d^2 r}{dt^2} - r \omega_n^2 \right] \hat{r} + T \left(2\omega_n \frac{dr}{dt} + r \frac{d\omega_n}{dt} \right) \hat{t} \\ &= T \left[\frac{d^2 r}{dt^2} - r |\nabla \times \vec{u}|^2 \right] \hat{r} + T \left(2|\nabla \times \vec{u}| \frac{dr}{dt} + r \frac{d|\nabla \times \vec{u}|}{dt} \right) \hat{t} \quad (14)\end{aligned}$$

The previous equation (14) reveals that the interaction between two matters in composed isolated system has existed along and across the radius at the same time, especially containing attractive force [14]. And the equation (14) can be changed into quantum form, if the vectors of \vec{r}

1 and \vec{u} have been changed into wave functions. As the simplest
 2 fundamental model of the motive matters' interaction in view of spin
 3 identified, it may be helpful to unite the four interactions in the field of
 4 space and velocity, which needs to be proved by experiments in future.

5 *Conclusions.*— The inertia and interaction have been studied for the
 6 matters with unchanged density in an isolated system. The inertia of spin
 7 with the angular velocity $\vec{\omega}_n = \frac{d\theta}{dt} \hat{n} = \nabla \times \vec{u}$ has been revealed and
 8 added into the Newton's the first law, for the matter of isolated system, by
 9 considering the point with non-zero volume but close to zero. The inertias
 10 of uniform angular spin and uniform rectilinear motion can be used to
 11 explain wave-particle duality and seismic waves. Further, the coupling
 12 isolated system with double matters has been chosen as the second
 13 fundamental model. The coupling characterizations inner two matters can
 14 be used to prove Newton's the second and third laws, to explain the
 15 structure of DNA, the shapes of Time Cone and topological sphere of
 16 moving trace. The force of two coupling matters has been expressed as a
 17 united formula $\vec{F} = T \left[\frac{d^2 r}{dt^2} - r |\nabla \times \vec{u}|^2 \right] \hat{r} + T \left(2 |\nabla \times \vec{u}| \frac{dr}{dt} + r \frac{d|\nabla \times \vec{u}|}{dt} \right) \hat{t}$,
 18 in which r and \vec{u} have the same origin of isolated system centre and
 19 can be instead of wave function to fit quantum mechanics, as shown in
 20 the literatures [15,16]. The interaction of coupling matters may be helpful
 21 to uniform the gravitation and electromagnetic force in future.

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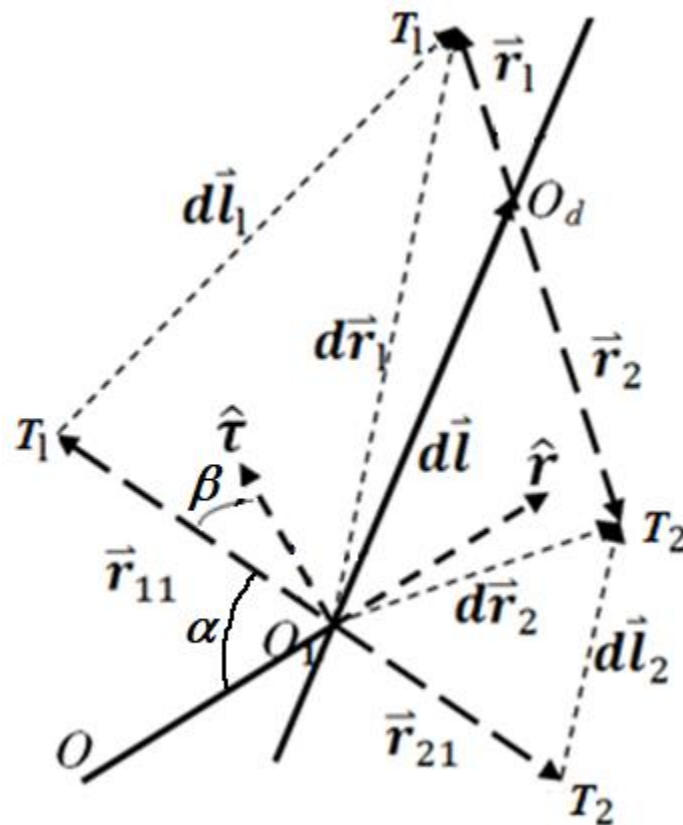


Figure 1